Galactic Structure and Dynamics

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AST 1420





Assignment 1 posted on the website!

Questions submitted

- Violent relaxation and phase mixing
- Relaxation time
- Virial theorem and GC application
- Collisionless Boltzmann equation
- Jeans equations
- Jeans theorem and distribution functions
- Applications:
 - Dwarf spheroidals
 - Local Group timing

Violent relaxation and phase mixing





Relaxation time

Relaxation time and collisional vs. collisionless systems

- triples, etc.?
 - encounters are very rare
 - only holds for the dynamics of 'field stars'
 - Use of the impulse approximation and ignoring very strong encounters:
 - expected velocity change.
 - This will be relevant when we discuss N-body simulations!

• Why are we able to ignore the dynamics of individual stars on each other? What about binaries,

• Basically, the relaxation-time argument is that the density is so low in galaxies that close

• Of course, this ignores binaries, stars born in clusters, and other stellar associations and

• For ~constant density medium, expected velocity-squared changes scale as db/b (b: impact parameter), so strong encounters only contribute a very small part of the total

Dense systems will through collisions evolve to a thermal state, similar to gas; galaxies do not



The viral theorem

The virial quantity

- $G = sum_i w_i x_i \cdot v_i$
- For wi = mi, G is twice the time derivative of the scalar moment of inertia (we'll
 use that later to derive the tensor version of the virial theorem)
- That G is conserved is an assumption about equilibrium
- If you don't set wi = mi, need to be careful to not break the assumption of equilibrium (so you can't be too specific about particle identity)
- wi =/= mi is used to create different mass estimators, but in practice need to test these with simulations

Milky Way mass from globular cluster kinematics

- Why does the self-gravitating case work so well?
 - as you combine these in a sensible way, won't be that far off!

• Estimator is still a ratio of squared velocities to inverse positions, so as long



Collisionless Boltzmann equation

Phase space density

- What is f(**x**,**v**) = f(**w**)?
 - Probability of seeing a body in vo density
 - If defined such that \int dx dv f(x, small volume
 - Always a probability of seeing a bas f(E) or f(E,L)

• Probability of seeing a body in volume dx dv, similar to spatial number

• If defined such that int dx dv f(x,v) = N, then it's the number of bodies in a

• Always a probability of seeing a body in this 6D volume, even when written

Jeans equations

Just a moment...

- What's a moment?
- Concept in probability theory, for probability • density p(x), moment is any function

 $int dx x^a p(x)$

[note that $\inf dx p(x) = 1$]

• For the Jeans equations, we consider moments of

 $p(\mathbf{v}) = f(\mathbf{x}, \mathbf{v}) / \operatorname{int} d\mathbf{v} f(\mathbf{x}, \mathbf{v}) = f(\mathbf{x}, \mathbf{v}) / \operatorname{nu}(\mathbf{x})$

[where nu(x) is the number density]





Integrating the CBE

$\frac{\partial f(\mathbf{q},\mathbf{p},t)}{\partial t} + \dot{\mathbf{q}} \frac{\partial f(\mathbf{q},\mathbf{p},t)}{\partial \mathbf{q}} + \dot{\mathbf{p}} \frac{\partial f(\mathbf{q},\mathbf{p},t)}{\partial \mathbf{p}} = 0,$

means that

$$\frac{\partial f(\mathbf{x}, \mathbf{v}, t)}{\partial t} + \dot{\mathbf{x}} \frac{\partial f(\mathbf{x}, \mathbf{v}, t)}{\partial t}$$

Integrating then gives

$$\int \mathrm{d}^3 \mathbf{v} \, \frac{\partial f}{\partial t} + \int \mathrm{d}^3 \mathbf{v} \, v_i \frac{\partial f}{\partial x_i} - \frac{\partial \Phi}{\partial x_i} \int \mathrm{d}^3 \mathbf{v} \, \frac{\partial f}{\partial v_i} = 0,$$

qdot and **pdot** are given by Hamilton's equations, for Cartesian coordinates this

 $\frac{(\mathbf{x},\mathbf{v},t)}{\partial \mathbf{x}} - \frac{\partial \Phi}{\partial \mathbf{x}} \frac{\partial f(\mathbf{x},\mathbf{v},t)}{\partial \mathbf{v}} = 0,$

Spherical Jeans equation and anisotropy

- Assuming zero rotation gives $\frac{\mathrm{d}(\nu \, v_r^2)}{\mathrm{d} r} + 2^{\frac{1}{2}}$
 - What if rotation is non-zero?
 - Assume some rotation and plug it in!
 - e.g., the globular-cluster system

$$\frac{\mathrm{d}(\nu \,\overline{v_r^2})}{\mathrm{d}r} + \nu \left(\frac{\mathrm{d}\Phi}{\mathrm{d}r} + \frac{2\overline{v_r^2} - \overline{v_\theta^2} - \overline{v_\phi^2}}{r} \right) = 0.$$

$$\frac{\beta}{r} \frac{1}{\nu v_r^2} = -\nu \frac{\mathrm{d}\Phi}{\mathrm{d}r} \,.$$

Dark matter halos are not believed to have much rotation, neither does.

Jeans theorem and distribution functions

Integrals of the motion

- Integrals of the motion I are quantities that depend on (\mathbf{x}, \mathbf{v}) and are conserved along the orbit
- They cannot explicitly depend on time
- Examples: energy in a static potential, angular momentum in a spherical potential
- Constants of the motion do depend on time and are not terribly useful
- Existence of integrals of the motion typically limits the orbit to a lower-dimensional subspace (e.g., energy, angular momentum). Those integrals are called *isolating*

The King profile

- Remember: King profile is a way to deal with the infinite nature of the isothermal sphere
- Do this by introducing a cut off in the distribution function

$$f(\mathcal{E}) = \frac{\rho_1}{(2\pi\,\sigma^2)^{3/2}}$$

- Where does this come from? Is there a good reason for this type of cut off?
 - Basically, it's one of the simplest ways to cut off the distribution function
 - Happens to fit elliptical galaxies and globular clusters well
 - There are other ways to cut it off that are less harsh

$$\left[\exp\left(\frac{\Psi-\frac{\nu^2}{2}}{\sigma^2}\right)-1\right]$$

Applications to masses of spherical systems

Dwarf spheroidals

- Why do they have so much dark matter?
- Question for galaxy formation! Multiple reasons:
 - Gas cooling in low-mass dark-matter halos is not efficient at all but the earliest times
 - Lowest mass dwarf spheroidals: gas is removed during reionization
 - Once galaxies fall into another galaxy, they cannot accrete more gas and thus their fuel for star formation is cut off
- All work together to increase the ratio of dark to baryonic matter in low mass galaxies

Local group timing

- Why do the Milky Way and M31 have to be in their first radial period?
- survive

• Because of their radial orbit, they pass through each other at the end of the first radial period and that would lead to a merger that the galaxies would not